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## Engineering of a Mo/SiN diffusion barrier to reduce the formation of MoS in CuZnSnS thin film solar cells

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**Engineering of a Mo/Si<sub>x</sub>N<sub>y</sub> Diffusion Barrier to Reduce the Formation of MoS<sub>2</sub> in Cu<sub>2</sub>ZnSnS<sub>4</sub> Thin Film Solar Cells**

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**Keywords:** CZTS, interface recombination, MoS<sub>2</sub>, TPV, solar cell and Si<sub>x</sub>N<sub>y</sub>.

**ABSTRACT:** The optimisation of the interface between back contact and absorber is one of the main challenges to improve the electrical behaviour and further enhance the efficiencies of Cu<sub>2</sub>ZnSn(S,Se)<sub>4</sub> (CZTS(e)) solar cell devices. In this work, Mo/Si<sub>x</sub>N<sub>y</sub> thin films with various film thicknesses were introduced as an interfacial layer to explore its influence on opto-electronic properties of the pure sulphide CZTS thin film solar cells. The Si<sub>x</sub>N<sub>y</sub> was deposited through plasma enhanced chemical vapour deposition (PECVD). The film thickness and stress of the Mo/Si<sub>x</sub>N<sub>y</sub> films were controlled to improve the adhesion of the CZTS layer and reduce the chances of cracking the deposited films. Energy dispersive X-Ray spectroscopy (EDS) mapping measurements performed directly on the cross-section of Mo/Si<sub>x</sub>N<sub>y</sub>/CZTS/Mo films indicate that the Si<sub>x</sub>N<sub>y</sub> intermediate layer can effectively inhibit the formation of a highly resistive MoS<sub>2</sub> layer and decomposition of CZTS at the CZTS/Molybdenum (Mo) interface region. A reduced efficiency was obtained with a Si<sub>x</sub>N<sub>y</sub> modified back contact compared with the devices without this layer. This could be due to the increased recombination and poor hole extraction stemming from the very low valance band maximum of Si<sub>x</sub>N<sub>y</sub> obtained from ultraviolet photoelectron spectroscopy (UPS) measurements. Temperature dependent current density-voltage (T-JV) and temperature dependent transient photovoltage (T-TPV) measurements were used to uncover insights into the internal recombination dynamics of the charge carriers.

**INTRODUCTION**

Pure sulphide CZTS thin film solar cells have been extensively studied in the last few years as an earth-abundant and environmentally-friendly alternative to well-established Cu(In,Ga)(Se,S)<sub>2</sub> (CIGS) technologies.<sup>1-3</sup> Pure sulphide CZTS is an attractive earth-abundant, environmental-friendly and stable light harvesting photovoltaic material. The bandgap of pure sulphide CZTS is around 1.5 eV, which is close to the optimal bandgap required for a single junction solar cell. However, the record power conversion efficiency of these pure sulphide CZTS solar cells and toxic selenium-containing CZT(S, Se) solar cells have been stagnant at around 9% and 12.6% for a few years, mainly due to a large V<sub>oc</sub> deficit.<sup>4-5</sup> Several factors has been proposed to explain this V<sub>oc</sub> deficit, including, an unfavourable thick MoS<sub>2</sub> layer at the back contact<sup>6</sup>, existence of charged defects associated with band tailing states<sup>7</sup>, improper band alignment at the CZTS/CdS heterojunction interface<sup>3</sup>. To form a

phase-pure CZTS(e) material with relatively large grains, a high temperature ( $>500\text{ }^{\circ}\text{C}$ ) sulfurization annealing step is essential and unavoidable. This would normally lead to the formation of  $\text{MoS}(\text{e})_2$  between Mo and CZTS(e). The functionality of this  $\text{MoS}(\text{e})_2$  is still debatable as it greatly depends on the thickness of the layer. A very thin layer of  $\text{MoS}(\text{e})_2$  may help to form a quasi-ohmic contact at the interface and improve the adhesion of CZTS(e) on Mo.<sup>8-9</sup> However, it is likely to form a quite thick layer in the absence of sulphur or selenium excess (low partial pressure) during high temperature annealing along with the decomposition of CZTS into secondary phases (e.g. SnS, ZnS), which could increase the series resistance and therefore is detrimental for solar cell performance.<sup>6, 8, 10</sup> To address this issue, an effective barrier layer was introduced to prevent the formation of thick  $\text{MoS}(\text{e})_2$ . To optimise the design of the back contact, several successful strategies have been successfully implemented to reduce the series resistance and increase the fill factor including Ag with good conductivity and additional p-type doping<sup>11</sup>,  $\text{ZnO}$ <sup>12-13</sup>,  $\text{MoO}_2$ <sup>14</sup> with good chemical stability;  $\text{TiN}$ <sup>8, 15-16</sup> and  $\text{TiB}_2$ <sup>9</sup> with both good chemical stability and electrical conductivity. More recently, bilayer structures  $\text{SiC}/\text{Mo}$ <sup>17</sup> and  $\text{Au}/\text{MoO}_3$ <sup>18</sup> were reported to mainly boost the open circuit voltage and hence improve the power conversion efficiency.

$\text{Si}_x\text{N}_y$  is a dielectric material with excellent chemical and thermal stability<sup>19</sup> and has been used as an antireflection coating (ARC) as well as a moisture barrier for crystalline silicon<sup>20</sup> and CIGS solar cells<sup>21</sup>. In this work, we employed a  $\text{Mo}/\text{Si}_x\text{N}_y$  diffusion barrier to avoid sulfurization of Mo and decomposition of CZTS. Due to the high resistivity of  $\text{Si}_x\text{N}_y$ , a further very thin layer of Mo (10-30 nm) was coated on top of  $\text{Si}_x\text{N}_y$ . We demonstrated the importance of controlling residual stress in back contact design to improve the adhesion of CZTS on the back surface. The influence of the  $\text{Mo}/\text{Si}_x\text{N}_y$  diffusion barrier on the materials quality and device performance of solar cells is presented.

## EXPERIMENTAL SECTION

To design a suitable back contact for solution processed CZTS, a series of surface treatments and coatings were performed on the Mo/soda-lime glass (SLG) substrates. Oxygen plasma treatment was used to remove any surface contamination and improve wetting of DMSO based C-Z-T-S solutions onto the Mo surface. All the precursor solution were doped with 0.14M NaCl. In this work, we employed a 10 nm-thick silicon nitride  $\text{Si}_x\text{N}_y$  layer using a PECVD method directly on top of the Mo back contact.  $\text{Si}_x\text{N}_y$  is a thermally and chemically stable material, which would serve as an appropriate barrier layer to reduce the sulfurization of Mo. DMSO-based solutions showed very poor wettability onto the  $\text{Si}_x\text{N}_y$  surface. Therefore in addition to the high resistivity of  $\text{Si}_x\text{N}_y$ , a 10-30 nm thick Mo capping layer were sputtered on top of the  $\text{Si}_x\text{N}_y$  layer. The additional Mo layer provide improved conductivity and adhesion of CZTS in comparison to the plain  $\text{Si}_x\text{N}_y$  layer. Furthermore, the effect of the  $\text{MoS}_2$  thickness formed between  $\text{Si}_x\text{N}_y$  and CZTS was investigated. The CZTS layers were deposited on these  $\text{Mo}/\text{Si}_x\text{N}_y$  modified back contacts by spin-coating of C-Z-T-S solution precursors and subsequently sulfurized in a rapid thermal processing (RTP, MTI Corporation) furnace at  $560\text{ }^{\circ}\text{C}$  for 20 mins. The solution preparation details are the same as those reported earlier.<sup>22</sup> The thickness of the CZTS absorbers was 1.2-1.3  $\mu\text{m}$ . After coating CZTS layers, a  $\sim 70\text{ nm}$  thick CdS layer was deposited by chemical bath deposition. Sputtered

intrinsic ZnO (~50 nm) and Sn:In<sub>2</sub>O<sub>3</sub> (~350 nm) served as a transparent conductive oxide. A Ni-Al metal grid was deposited using thermal evaporation to improve the conductivity of the device.

The morphology of films was studied using a JEOL-JSM-7800F field emission scanning electron microscope (5 kV acceleration voltage, a working distance of 10 mm and a magnification of x 25, 000). Energy dispersive X-ray spectroscopy (EDX) mapping was used to deduce the element distribution using 20 kV acceleration voltage. The stress measurements were made with a Bruker D8 Discover X-ray diffraction system with a 0.07 ° step size, at a time of 7 seconds per step. The undertaken scans were over the full 0-0.9  $\sin^2(\psi)$  in both positive and negative  $\psi$  tilts to confirm the absence of shear stress. Secondary ion mass spectrometry (SIMS) depth profiles were obtained using an in-house instrument constructed at the Interface Analysis Centre, University of Bristol, comprising of an electronically variable aperture type gallium ion gun (FEI SD gallium LMIS EVA focusing column) fitted to a double focusing magnetic sector mass analyser (Vacuum Generators model 7035). Surface spectra were obtained by scanning a 25 keV gallium ion beam over an area of the sample while scanning the magnet to accept secondary ions in the range 0-125 Daltons, with a step size of 0.05 Daltons and a dwell time of 100 ms per step. Hence a total time of 200 s to obtain the spectrum. Depth profiles were obtained by scanning a 3 nA gallium ion beam over an area of 43  $\mu\text{m}$  x 43 $\mu\text{m}$  or 65  $\mu\text{m}$  x 65 $\mu\text{m}$  and monitoring the species of interest. Sodium, silicon, copper, zinc, molybdenum and tin signals were counted with dwell times of 1s each, cycling through the elements for a total period of 20 minutes. The system vacuum during operation was  $2 \times 10^{-8}$  mbar. The instrument was calibrated using values of 68.93 for the backscattered Ga<sup>+</sup> ions. The instrument control software was 'Pisces' running under the Windows operating system.<sup>23</sup> UPS was performed alongside XPS in the Kratos Axis Supra, utilising the He(I) line with a pass energy of 10 eV, step size of 0.025 eV and a 65 ms dwell. A 110  $\mu\text{m}$  aperture was utilised in order to prevent saturation of the detector. The secondary electron cut-off and valence band edges were estimated using the Step Up and Step Down edge backgrounds within CasaXPS. In each case the measured value is the intersection of the step with the background electron counts. The current-density-voltage (J-V) curves for the solar cell devices were measured under simulated AM 1.5G spectrum and 100 mW/cm<sup>2</sup> (1 sun) illumination. The external quantum efficiency (EQE) measurements were performed in AC mode with a chopping frequency of 67 Hz using a QEX10 system (PV Measurements) calibrated with a NIST-certified Si photodiode. Transient photovoltage measurements were performed using a commercially available transient measurement system (Automatic Research GmbH). This system uses a 635 nm red laser diode driven by a waveform generator (Keysight 33500B). The laser pulse length was 100 ns. Background illumination was provided by a white LED with its intensity calibrated to generate the same device photocurrent as measured using the solar simulator. This intensity is referred to as 1 Sun equivalent. Transient responses were captured by a digital storage oscilloscope (Keysight DSOX2024A), the number of sample averages being adjusted to optimise signal noise and measurement time. The device under test is assumed to be held at open-circuit by the 1 M $\Omega$  oscilloscope input. TPV decays were fitted using a single exponential function. Temperature

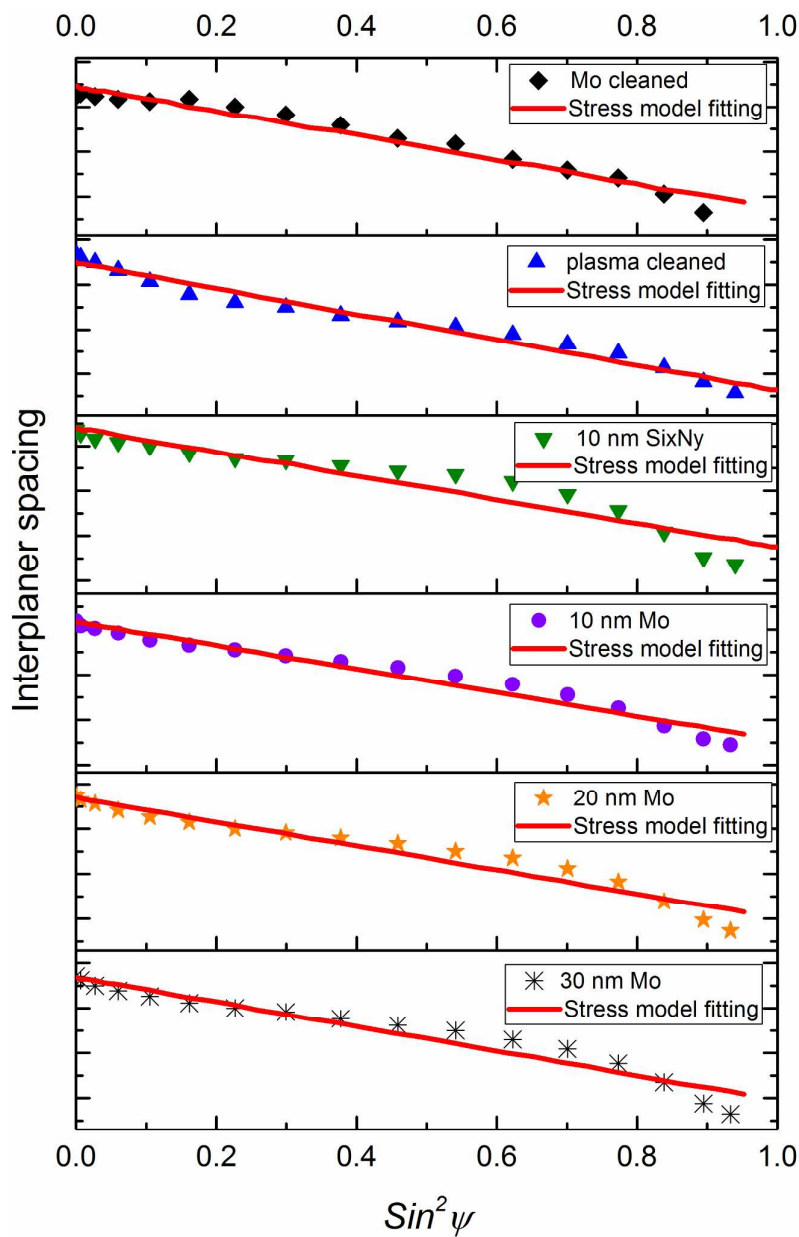
dependent J-V and TPV measurements were performed using a temperature-controlled chamber (Linkam) under liquid nitrogen cooling.

## RESULTS AND DISCUSSION

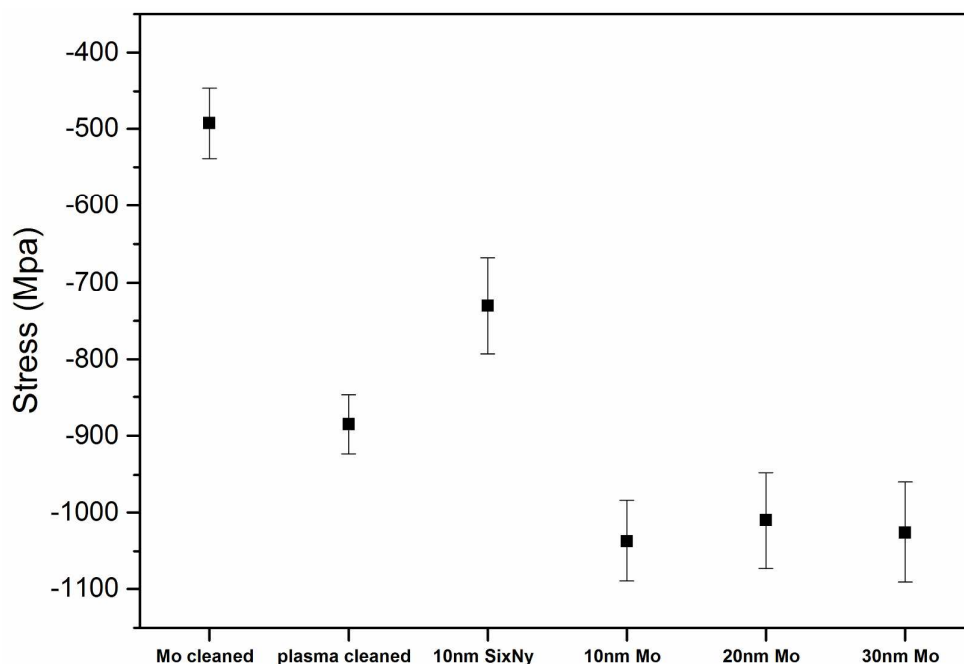
Since  $\text{Si}_x\text{N}_y$  and CZTS were prepared at 400 °C and 560 °C respectively, thermal effects would further enhance the stress variation which was indicated by CZTS film delamination in a previous publication<sup>24</sup>. Therefore, proper control on the stress is important for the back contact design. To determine the residual stress in multiple layer coatings, XRD was used to measure the peak shift on the diffraction pattern associated with inter-planar spacing  $d$  with relation to  $\psi$  angle. Strain causing elongation or contraction of the crystal lattice will induce the change in  $d$ , which could be obtained on the diffraction pattern. In this work, the  $\sin^2\psi$  plot was used to derive the residual stress  $\sigma$  using Equation 1:

$$\sigma = \frac{E}{(1+\nu)\sin^2\psi} \left( \frac{d_\psi - d_n}{d_n} \right) \quad \text{Equation 1}$$

where the inter-planar spacing of planes normal to the surface  $d_n$ , the inter-planar spacing of planes at an angle  $\psi$  to the surface  $d_\psi$ , Elastic modulus  $E$  and Poisson's ratio.<sup>25</sup> A number of XRD measurements were made at different psi ( $\psi$ ) tilts for the Mo cleaned, plasma treated,  $\text{Si}_x\text{N}_y$ ,  $\text{Si}_x\text{N}_y$  coating with 10 nm, 20 nm and 30 nm Mo-capping as shown Figure 1. The calculated stress values extracted from Figure 1 for the individual modifications of the back contact were plotted in Figure 2. The bottom Mo film was cleaned with DI water and IPA, the residual stress is estimated as  $-492.4 \pm 46.5$  Mpa, which was used as a baseline for all the stress measurements. After 10 minutes  $\text{O}_2$  plasma cleaning on the Mo surface, a dramatic increase in the compressive stress is derived as  $-885.1 \pm 37.8$  Mpa. As we observed a better wetting after  $\text{O}_2$  plasma cleaning, this dramatic increase of the compressive stress could be a good indication for improved film adhesion. However, a pronounced reduction of the compressive stress can be calculated as  $-730.4 \pm 62.6$  Mpa after  $\text{Si}_x\text{N}_y$  coating. Changes in the mechanical property of the materials or the Mo/ $\text{Si}_x\text{N}_y$  stacked layers may contribute to the delamination of sulphurised CZTS films and poor wetting of the C-Z-T-S solution on  $\text{Si}_x\text{N}_y$  surface. A significant increase in compressive stress after sputtering of the 10-30 nm Mo capping layer is estimated in the range of -1040 to -1000 Mpa. Improved wetting of the C-Z-T-S solution and CZTS film adhesion onto the  $\text{Si}_x\text{N}_y$  surface was obtained whilst in this stress range.



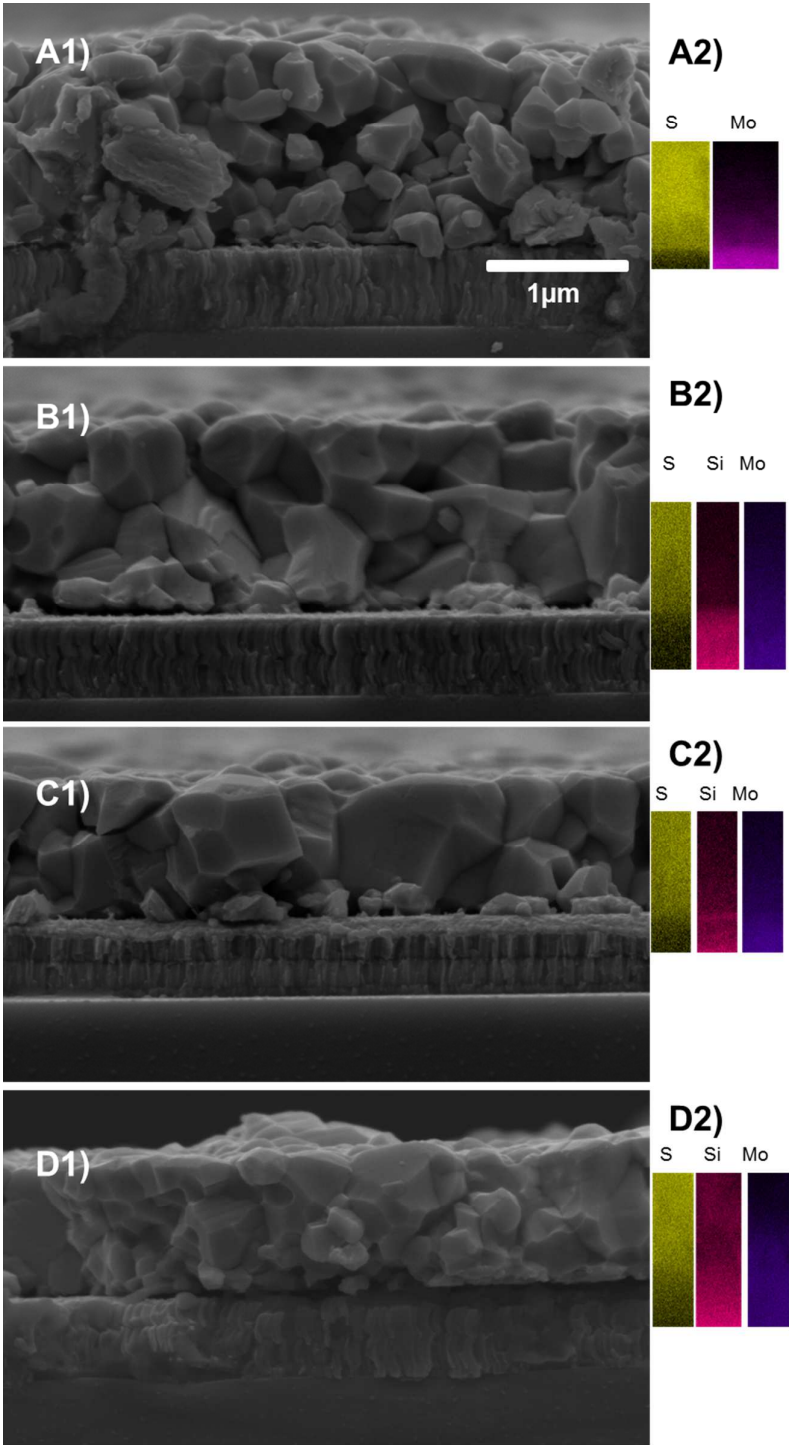
**Figure 1.**  $\sin^2\psi$  plot of experimental and modelled results of various back contact configurations (untreated Mo, plasma cleaned Mo, 10nm-Si<sub>x</sub>N<sub>y</sub>/Mo/SLG, 10nm-Mo/10nm-Si<sub>x</sub>N<sub>y</sub>/Mo/SLG, 20nm-Mo/10nm-Si<sub>x</sub>N<sub>y</sub>/Mo/SLG, 30nm-Mo/10nm-Si<sub>x</sub>N<sub>y</sub>/Mo/SLG)



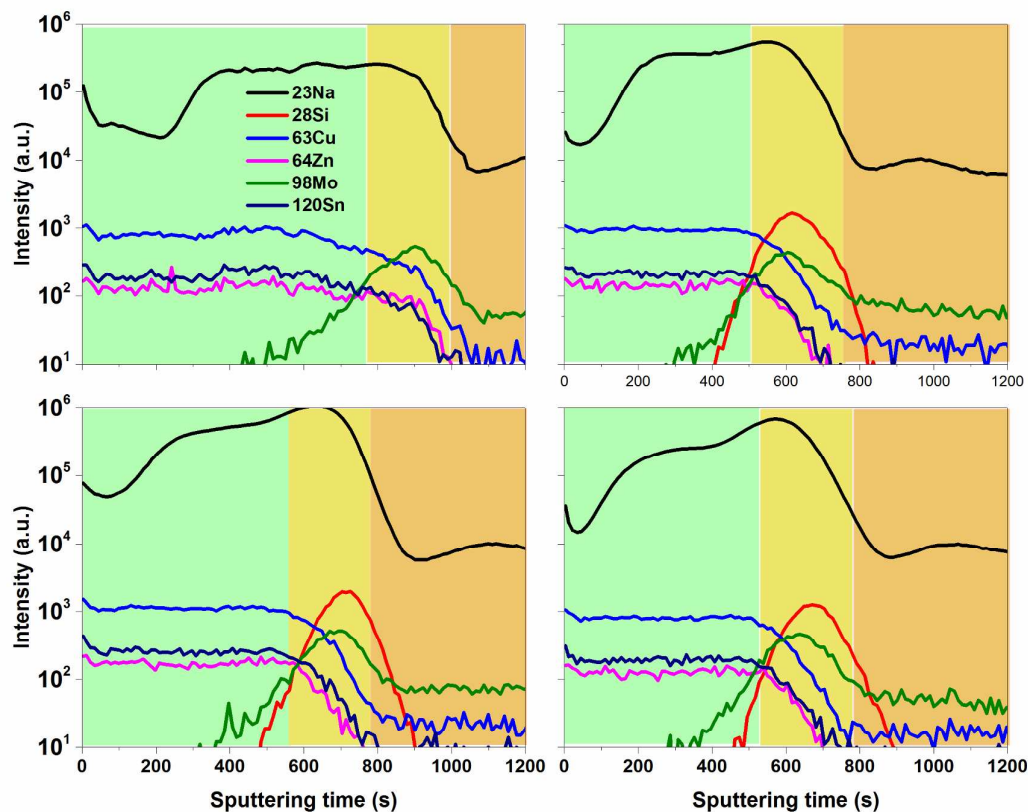
**Figure 2.** Residual stress of various back contact configurations (untreated Mo, plasma cleaned Mo, 10nm-Si<sub>3</sub>N<sub>4</sub>/Mo/SLG, 10nm-Mo/10nm-Si<sub>3</sub>N<sub>4</sub>/Mo/SLG, 20nm-Mo/10nm-Si<sub>3</sub>N<sub>4</sub>/Mo/SLG, 30nm-Mo/10nm-Si<sub>3</sub>N<sub>4</sub>/Mo/SLG)

After tuning the individual deposition parameters for Si<sub>3</sub>N<sub>4</sub> and Mo capping layers, the stress mismatch between the CZTS and back contacts was improved. To investigate the effectiveness of Si<sub>3</sub>N<sub>4</sub> and Mo capping layers as a sulphurisation barrier, cross-section SEM and EDS mapping were used to characterise CZTS samples with and without barrier layers (Figure 3). CZTS with good adhesion were obtained on Mo-capping/Si<sub>3</sub>N<sub>4</sub> as shown in Figure 3A1-3D1. CZTS films grown directly onto the Si<sub>3</sub>N<sub>4</sub> surface were non-uniform due to poor wetting and adhesion. The CZTS films with the Mo-capping/Si<sub>3</sub>N<sub>4</sub> barrier layers (Figure 3B1, 3C1 and 3D1) show a slightly bigger grain size than those without barrier layers (Figure 3A1). No significant differences in the film thickness and film compactness were found in these films. Figure 3A2 shows a sulphur distribution into the underlying Mo layer which could be an indication of MoS<sub>2</sub> formation. For the samples with barrier layers, a clear interface was created between sulphur and silicon and very little sulphur was detected within the Mo region (Figure 3B2, 3C2 and 3D2). The Mo-capping/Si<sub>3</sub>N<sub>4</sub> bilayer is shown to reduce or avoid the formation of MoS<sub>2</sub> at the back contact.



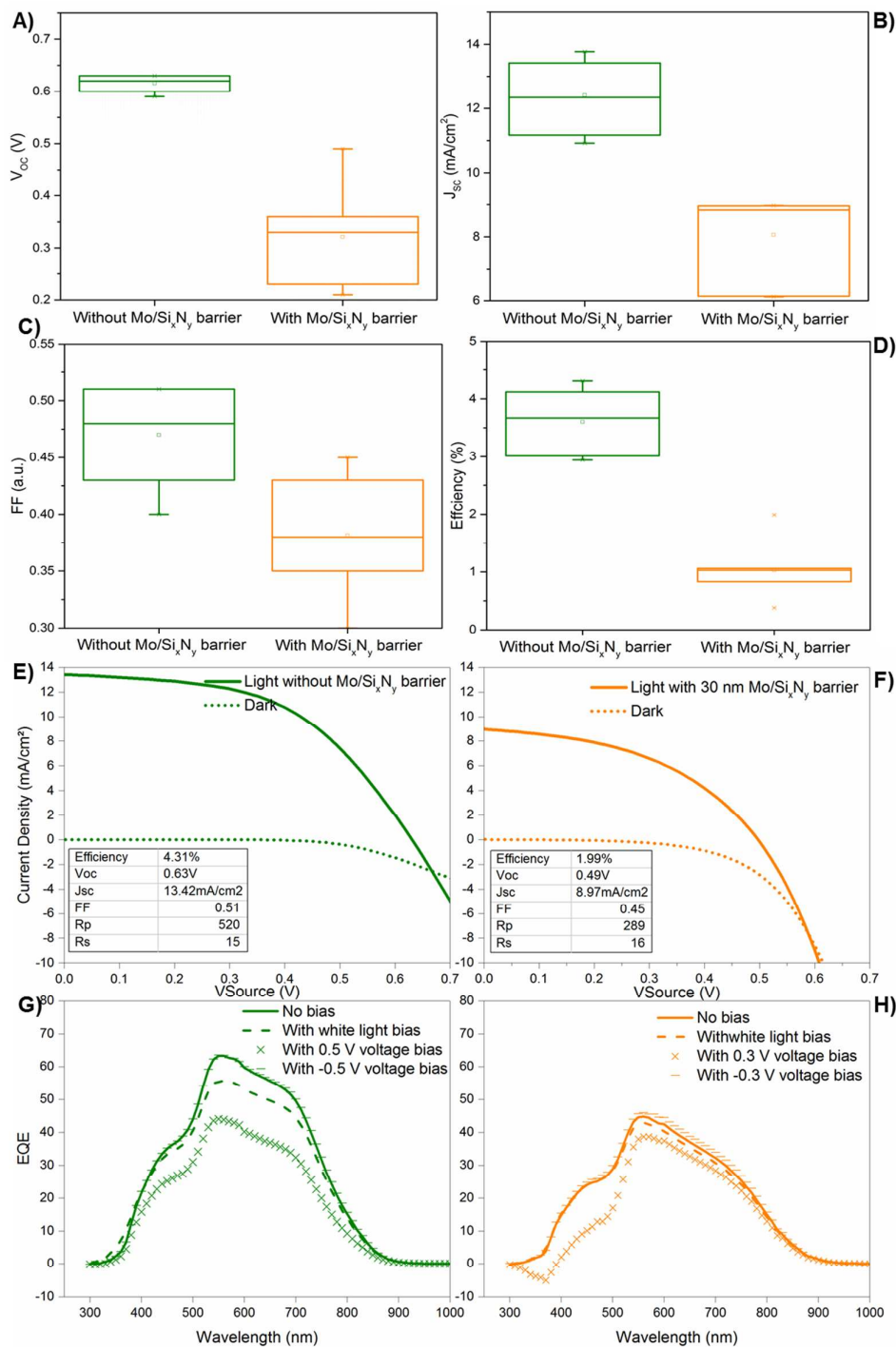


**Figure 3.** Cross-section SEM images of CZTS samples on Mo/SLG with or without barrier layers: A1) CZTS/Mo/SLG, A2) EDS mapping of a select area from SEM image A1); B1) CZTS/10nm-Mo/10nm-Si<sub>x</sub>N<sub>y</sub>/Mo/SLG, B2) EDS mapping of a select area from SEM image B1); C1) CZTS/20nm-Mo/10nm-Si<sub>x</sub>N<sub>y</sub>/Mo/SLG, C2) EDS mapping of a select area from SEM image C1); D1) CZTS/10nm-Mo/30nm-Si<sub>x</sub>N<sub>y</sub>/Mo/SLG, D2) EDS mapping of a select area from SEM image D1). The selected the elements in EDS mapping are labelled as yellow for S, purple for Mo and red for Si.



**Figure 4.** SIMS elemental depth profiles of CZTS samples on Mo/SLG with or without barrier layers: A) CZTS/Mo/SLG, B) CZTS/10nm-Mo/10nm-Si<sub>x</sub>N<sub>y</sub>/Mo/SLG, C) CZTS/20nm-Mo/10nm-Si<sub>x</sub>N<sub>y</sub>/Mo/SLG, D) CZTS/10nm-Mo/30nm-Si<sub>x</sub>N<sub>y</sub>/Mo/SLG. CZTS, MoS<sub>2</sub>/Si<sub>x</sub>N<sub>y</sub> and Mo were showed with green, yellow and orange backgrounds, respectively.

From the SIMS depth profiles (Figure 4), there isn't significant difference on the depth profiles of Na due to intentionally doped all CZTS precursors with 0.14M NaCl. However, there is an apparent reduction of the Mo diffusion length after Mo/Si<sub>x</sub>N<sub>y</sub> barrier was added in Sample B-D. The slight variation of Mo intensity could be due to the varied thickness of MoS<sub>2</sub> formed towards the Si<sub>x</sub>N<sub>y</sub> surface. Further investigation with TEM/EDS would be required to analyse the interfaces in nm/sub-nm level.



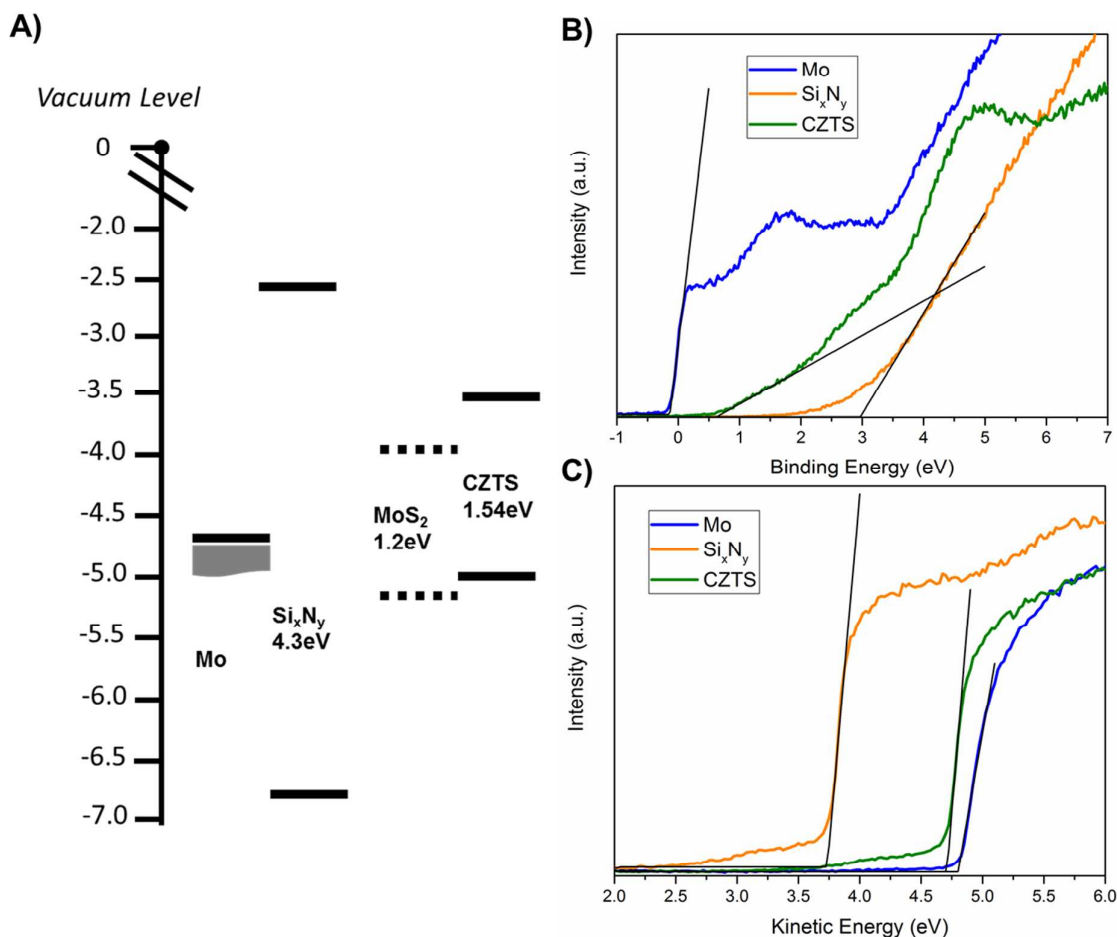
**Figure 5.** (A-D) Box charts of J-V characteristics of CZTS devices with and without Mo/Si<sub>x</sub>N<sub>y</sub> barrier layer. J-V curves for the best-performing CZTS solar cells fabricated E) without Mo/Si<sub>x</sub>N<sub>y</sub> barrier layers and F) with Mo/Si<sub>x</sub>N<sub>y</sub> barrier layers. Solid lines represent J-V curves under illumination and dotted lines represent J-V curves in the dark condition. Basic device parameters were inserted in the graphs. EQE characteristics for the best-performing CZTS solar cells fabricated G) without Mo/Si<sub>x</sub>N<sub>y</sub> barrier layers and H) with Mo/Si<sub>x</sub>N<sub>y</sub> barrier layers under various bias conditions.

The 10 nm-, 20 nm- and 30 nm-Mo-capping/ $\text{Si}_x\text{N}_y$ /Mo/SLG stacks were used to fabricate CZTS devices and compared against the reference CZTS device directly deposited onto Mo/SLG. Unfortunately, most devices with Mo/ $\text{Si}_x\text{N}_y$  back contact showed a significant reduction in open-circuit voltage and short-circuit current as shown in Figure 5A-D. The highest efficiency of 1.99% was achieved with Mo-30 nm/ $\text{Si}_x\text{N}_y$  barrier layer (Figure 5F) with significant reduction in  $V_{oc}$  and  $J_{sc}$  along with some decrease in FF as compared to solar cells without barrier layers (Figure 5E). The cross-over behaviour between the light and dark  $J$ - $V$  curve could be an indication of the presence of a blocking back contact at the interface between the Mo and CZTS.<sup>26</sup>

Light-bias-dependent and voltage-bias-dependent EQE were performed to investigate the efficiency loss mechanisms of fabricated CZTS solar cells as shown in Figure 5G and 5H. The minimum bandgap was estimated to be 1.58 eV and 1.54 eV from the absorption onset of EQE curves for the fabricated CZTS solar cells with and without barrier layers, respectively. Under white light bias (approximately 1 sun), no pronounced difference in the blue EQEs were observed for both samples as compared to that in the dark. This could be because the monochromatic blue probing beam has lower intensity than the blue part of the white bias light.<sup>27</sup> Thus it is hard to measure an increased EQE signal from charge carriers generated in the CdS region. However, noticeable drops in the wavelength range of 500 to 750 nm under illumination were seen compared to those in the dark. This reduction could be caused by the extra photogenerated electrons in the conduction band, which reduces the effective work function of CZTS and leads to lower collection efficiency.<sup>28</sup>

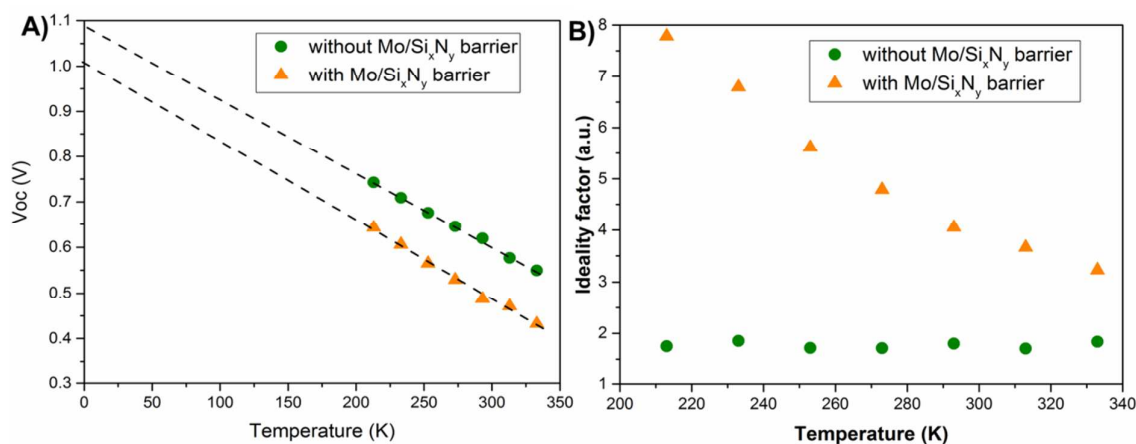
Voltage-bias-dependent EQE was used to create a forward bias and a reverse bias on the two devices. To avoid any damage caused by the bias voltage, the bias voltages ( $\pm 0.5$  V and  $\pm 0.3$  V for the solar cell with and without Mo/ $\text{Si}_x\text{N}_y$  barrier layers) that were chosen are slightly lower than  $V_{oc}$  of two devices (Figure 5G and 5H). Compared to the EQE under no bias condition (in the dark), the EQE increases marginally at longer wavelength with applied -0.5 V (Figure 5G) and -0.3 V (Figure 5H) reverse voltage biases. This is mainly contributed by the increased depletion width in CZTS under reverse voltage bias which helps the collection efficiency of red photons.<sup>27</sup> This observation also indicates high recombination losses and low minority carrier lifetime within the p-type CZTS absorber layers.<sup>28</sup> This low lifetime could stem from a high defect concentration in the absorber layer or high recombination losses at the back contact or at the front interface.<sup>26</sup> The EQE decreases across the whole spectrum range with applied +0.5 V (Figure 5G) and +0.3 V (Figure 5H) forward voltage biases compared with the zero voltage bias. This could be because the depletion width in CZTS reduces significantly with forward voltage bias, hence a much lower carrier collection of photogenerated minority carriers.<sup>29</sup> It is notable that the negative EQE values were obtained at the wavelength between 300 nm and 400 nm under reverse bias of +0.3 V. Although applied forward voltage is still lower than  $V_{oc}$ , the dark current caused by majority carrier electrons in n-type CdS is larger than the photogenerated current, therefore an overall negative EQE is obtained. The reduced built-in electrical field also reduces the number of drift electrons moving towards the p-n junction.<sup>28</sup> The device with Mo/ $\text{Si}_x\text{N}_y$  barrier layers is

more sensitive to the reduction of depletion width, which could be partially caused by accumulation of holes at the Mo/Si<sub>x</sub>N<sub>y</sub> barrier layers.



**Figure 6.** A) Band diagram of the effective band edge states of the CZTS and modified back contact layers in solid lines. B) the Fermi level and C) work functions are measured by UPS which were used to calculate the ionisation potential (or valence band maximum) of each film. The band gap of the CZTS and Si<sub>x</sub>N<sub>y</sub> were obtained using EQE and UV-VIS, respectively. Band edge states of the MoS<sub>2</sub> were adapted from reference <sup>30</sup>, which could not be detected by UPS in this work.

To investigate differences in energetics after applying the Mo/Si<sub>x</sub>N<sub>y</sub> barrier layers, we employed ultraviolet photoelectron spectroscopy (UPS). Ionisation potentials (IP or valence band maximum) of  $(6.8 \pm 0.1)$  eV and  $(5.1 \pm 0.1)$  eV were measured for Si<sub>x</sub>N<sub>y</sub> and CZTS films, respectively, as shown in Figure 6. Mo thin films were found to have a work function of  $(4.7 \pm 0.1)$  eV. A 2 eV greater barrier height was introduced by 30 nm-thick Mo/Si<sub>x</sub>N<sub>y</sub> layers compared to unmodified CZTS. This difference would form a blocking back contact at the interface between the Mo and CZTS, which can suppress the majority carrier (hole) transport.



**Figure 7.** A) Temperature dependency of open-circuit voltage ( $V_{oc}$ ) for CZTS solar cells without Mo/Si<sub>x</sub>N<sub>y</sub> barrier layers (solid circles) and with Mo/Si<sub>x</sub>N<sub>y</sub> barrier layers (solid triangles), the linear extrapolation to  $T = 0$  K in dashed lines. Data for both devices have 0 K intercepts that do not reach the band gap value. B) The correlated ideality factors versus temperature data for without Mo/Si<sub>x</sub>N<sub>y</sub> barrier layers (solid circles) and with Mo/Si<sub>x</sub>N<sub>y</sub> barrier layers (solid triangles).

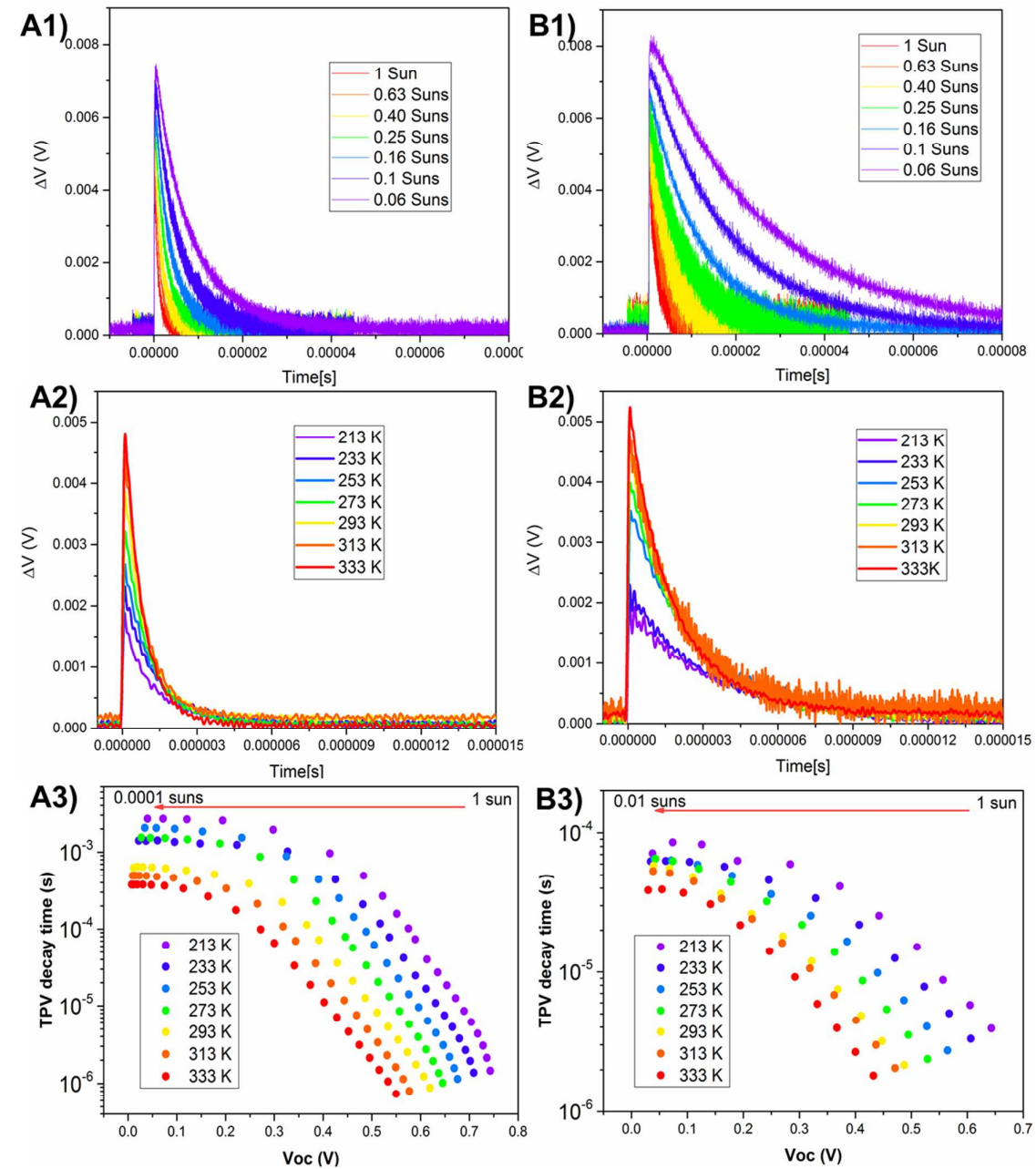
To further investigate the deficiency of the performance of the CZTS device with a Mo/Si<sub>x</sub>N<sub>y</sub> barrier as compared with the CZTS device without this barrier, temperature dependent  $V_{oc}$  data has been collected to determine the dominant recombination processes. The relationship between  $V_{oc}$  and temperature is according to <sup>31</sup> and as shown below:

$$V_{OC} = \frac{E_a}{q} - \frac{AkT}{q} \ln \frac{J_{00}}{J_L} \quad \text{Equation 2}$$

where  $E_a$ ,  $A$ ,  $J_{00}$ ,  $J_L$ ,  $q$ ,  $k$  and  $T$  are the activation energy for dominant recombination mechanism, diode ideality factor, reverse saturation current prefactor, photocurrent, electron charge, Boltzmann constant and the temperature, respectively. The  $V_{oc}$  vs  $T$  plots are shown in Figure 7A for the CZTS devices with and without Mo/Si<sub>x</sub>N<sub>y</sub> barrier layers, which yield 0 K intercept of the linear extrapolation of  $E_a/q$  values of 1.08 and 1.01 eV, respectively. Those  $E_a/q$  values are low compared to their respective band gap values (1.58 and 1.54 eV) which obtained from EQE measurement. This result indicates that the main recombination mechanism in these CZTS cells is dominated by interface recombination.<sup>26</sup> The more pronounced reduction on  $E_a/q$  value of the CZTS cell with the Mo/Si<sub>x</sub>N<sub>y</sub> barrier layers could be due to increased interface recombination occurring at the back contact with improper band alignment (Figure 7A). The efficiency for both devices collapses at low temperature due to increasing series resistance and the possible charge carrier freeze-out effect at low temperature as agreed with other publications.<sup>26, 32</sup> In Figure 7B, the T-dependent diode ideality factors for both CZTS solar cells with and without Mo/Si<sub>x</sub>N<sub>y</sub> barrier layers were plotted. The diode ideality factors for the CZTS solar cell without the Mo/Si<sub>x</sub>N<sub>y</sub> barrier are between 1 and 2, showing a marginal variation in the measured temperature range, which could be an indication of the deep defects acting as dominant trap states or so called Shockley-Read-Hall recombination via band tails.<sup>27, 33</sup> However, a significant increase of



diode ideality factor ( $A>2$ ) was observed for CZTS solar cells with the Mo/Si<sub>x</sub>N<sub>y</sub> barrier at low temperatures. The origin of  $A>2$  is not only by considering a single recombination route, but involves a multistep recombination process via a series of trap states distributed within the solar cell device structure.<sup>34</sup> This would indicate strong back contact recombination which could lower  $V_{oc}$ .<sup>26, 34-35</sup>



**Figure 8.** The representative TPV decays analysis of CZTS solar cells without Mo/Si<sub>x</sub>N<sub>y</sub> barrier layers (A) and with Mo/Si<sub>x</sub>N<sub>y</sub> barrier layers (B). A1) and B1) The representative TPV decay as a function of white-light bias in the range of 0.06 - 1 suns at 293 K. A2) and B2) Temperature-dependent TPV decay under 1 sun white-light bias at the temperature range of 213 – 333 K. A3) and B3) TPV decay time characteristics of CZTS solar cells under WLB range of 0.0001 – 1 suns and temperature range of 213 – 333 K.

To understand how the charge transport behaviour at the CZTS/back contact interface differs between the device types, we conducted temperature-dependent TPV measurements. For TPV measurements, a light pulse is applied to a device held at open-circuit condition under a white-light bias (WLB), and the resulting photovoltage changes created by the laser pulse ( $\Delta V$ ) is measured over time. From this, we measured the decay time of the open-circuit voltage as a function of the intensity of WLB, shown in Figure 8A1 and 8B1. TPV decay measurement gives insight into the internal recombination dynamics of the charge carriers. The decay time represents the internal extraction time to the contacts of the charge carriers before they recombine or are trapped in the defects.<sup>36</sup> The longer decay time measured, the higher chance the photo-generated charge carriers can be extracted. The decay times of charge carriers for the CZTS solar cells with and without Mo/Si<sub>x</sub>N<sub>y</sub> barrier layers are estimated on the order of  $\sim \mu\text{s}$  in the WLB range of 0.06 – 1 suns. With increasing WLB,  $V_{oc}$  of both CZTS devices increases. This could be caused by shallow defects that were fulfilled at higher excitation densities.<sup>37</sup> By increasing the WLB, there are more photo-generated minority charge carriers (electrons) in the CZTS absorber which could also increase the internal recombination with holes, reducing the extraction time. The variation on decay time of CZTS with Mo/Si<sub>x</sub>N<sub>y</sub> barrier layers using different light intensities is more significant than the CZTS device without the Mo/Si<sub>x</sub>N<sub>y</sub> barrier layers. This would be a sign of stronger recombination caused by this modified back contact. To investigate the thermally activated charge transport, temperature dependent TPV decay for the CZTS solar cells with and without Mo/Si<sub>x</sub>N<sub>y</sub> barrier layers at different temperature under 1 sun WLB were shown in Figure 8A2 and 8B2. The decay time of both CZTS devices increases with decreasing temperature. The longer extraction time could be due to reduced thermal energy to assist the charge carrier to jump spatially over fluctuated electrostatic potential<sup>7</sup> and travel across the CZTS device. A more pronounced variation on  $\Delta V$  and decay time for CZTS device with Mo/Si<sub>x</sub>N<sub>y</sub> barrier layers at different temperature as compared to the CZTS device without the barrier layers was obtained. To find out the combined effects of the intensity of WLB and temperature on TPV decay time, TPV decay time versus  $V_{oc}$  at different conditions were shown in Figure 8A3 and 8B3. The data were recorded until  $V_{oc} \sim 0$  at lowest light levels of WLB. We could still obtain  $V_{oc}$  at even 0.0001 suns for the CZTS solar cell device without Mo/Si<sub>x</sub>N<sub>y</sub> barrier layers, which is nearly 100 times lower light intensity than for CZTS device with Mo/Si<sub>x</sub>N<sub>y</sub> barrier layers. In other words, the increased interfacial recombination in the cell with Mo/SiN causes a significant decrease in  $V_{oc}$  at low light intensities. The TPV signal decays 100 times faster for a given  $V_{oc}$  after introducing the Mo/Si<sub>x</sub>N<sub>y</sub> barrier into device structure. These results indicate that the recombination within the CZTS device with Mo/Si<sub>x</sub>N<sub>y</sub> barrier layers may not only account for trapping and detrapping process<sup>38</sup>, but can also be attributed to the increased interface recombination between accumulated holes near the high barrier layer and free electrons in CZTS. This is in good agreement with the hole blocking effect observed from temperature dependent  $J-V$  measurement.

## CONCLUSIONS

In conclusion, we have engineered Mo/Si<sub>x</sub>N<sub>y</sub> barrier layers into a back contact structure for CZTS solar cells to effectively reduce the formation of the detrimental MoS<sub>2</sub> layer, which



was revealed by EDS mapping on the cross-section of CZTS samples with and without Mo/Si<sub>x</sub>N<sub>y</sub> barrier layers. We also demonstrated that stress measurements can be used to monitor the dramatic changes after the addition of every back contact layers, providing empirical guidance to control the adhesion of these films. A power conversion efficiency drop of 2.32% was observed after introducing the modified back contact. This reduction of efficiency is mainly due to the increased recombination towards back contact, which stem from the misalignment of valence band maximum between Si<sub>x</sub>N<sub>y</sub> and CZTS obtained by UPS. We have also demonstrated a combination characterisation methods to reveal the recombination information within the CZTS solar cells, particularly temperature dependent J-V and TPV showed an increased in interface recombination and poor hole extraction due to misalignment of band edges between Mo/Si<sub>x</sub>N<sub>y</sub>/CZTS region.

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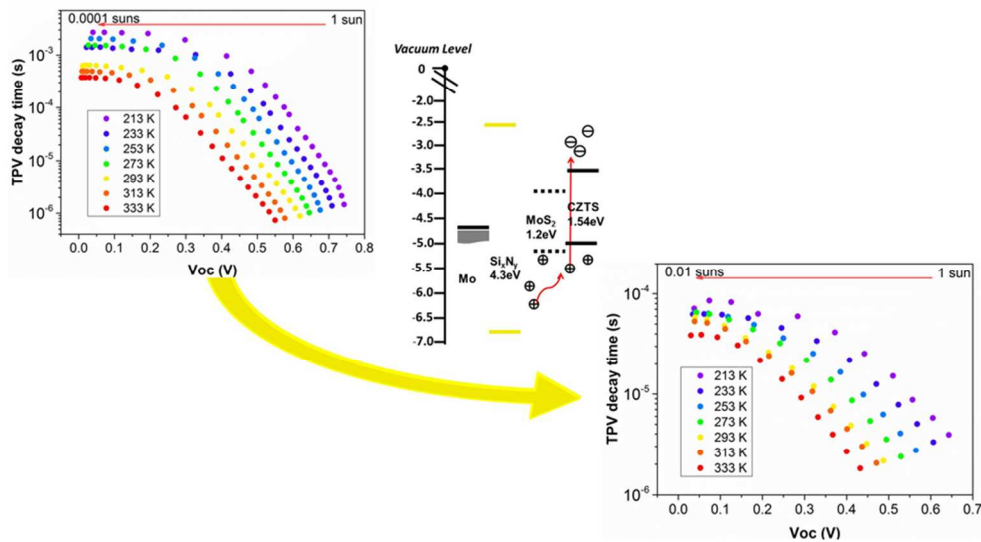
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The figure consists of two main parts. The top part is a schematic diagram of the MoS<sub>2</sub>/ZnTe heterostructure. It shows a cross-section of the device with MoS<sub>2</sub> (4.36 eV) and ZnTe (5.46 eV) layers. A vacuum level diagram is shown above the schematic, indicating the energy levels of the materials. The bottom part is a plot of TPD decay time (s) versus Voc (V) for various temperatures (213 K, 233 K, 253 K, 273 K, 293 K, 313 K, 333 K). The plot shows that the TPD decay time increases with Voc and decreases with temperature. A large yellow arrow points from the schematic to the spectra.



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